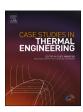
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# Influence of space conjugate temperature varying non-uniform heat sink/source on hydromagnetic slip water-EG (50:50) nanofluid

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#### ABSTRACT

Significant cooling is an essential requirement in the field of aeronautical and bio-medical engineering, nuclear reactor system, solar collectors, and in development of electronic chips etc. Involving rotating conical geometries. In view of this, flow and heat transfer over conical geometries subject to different constraints of motion are highly needed. Implementation of magnetic field subject to flow pattern controls the fluid motion thereby imparting better cooling. Consideration of nanofluid instead of regular fluid yields prominent cooling of the associated surface. Such relevance has motivated the authors to work on magnetohydrodynamics and heat transfer investigation of water-EG (50:50) mixture based  $Al_2O_3$  and  $Fe_3O_4$  past heated and rotating down-pointing upright cone subject to impact of space and temperature varying non-uniform heat source or sink. Numerical solution of dimensionless governing equations is accomplished by implementing Runge-Kutta method. The findings indicate that swirl and axial velocities peter out with rise in magnetic parameter while both exhibit opposite impact in response to slip parameter. Temperature profiles upgrade due to amplification of space and temperature dependent parameters. Skin friction and heat transportation upsurge with growth of solid volume fraction of nanoparticle.

#### 1. Introduction

Customizing the execution of industrial tools in the field of heat transfer (HT) is very challenging one. Early stage of research is mainly aimed using base fluids for HT phenomenon, but later on, to overcome this, Choi and Eastman [1] came up with an innovative type of fluids called nanofluids (NFs), which is a boon in enhancing the thermal conductivity. In recent years significant methods are developed to modify thermal performance of fluids by putting nano sized solid particles like  $Al_2O_3$ , Cu, Ag etc. in fluids, which are poor in HT. Hakeem et al. [2] conducted a comparative study with magnetized/non-magnetized nanoparticles. Soret-Dufour impacts

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on magnetized Casson NF flow across an inclined plane were perused by Lorenzini et al. [3]. They found that the solutal and thermal distributions improve and velocity slows down with Casson parameter. The impact of Brownian movement and thermophoresis on chemically-reacting magnetized Hiemenz stream across a nonlinear expandable plate was investigated by Ibrahim et al. [4]. Impacts of slip on magnetic Casson nanofluid flow over a permeable stretched sheet were perused by Kumar et al. [5]. 3D NF stirring with a non-uniform heat sink/source was studied by Thumma et al. [6]. They discovered that an upsurge in the stretching ratio parameter had an impact on the axial velocity profile. Sheri and Thumma [7] investigated the thermal and solutal effects on natural convection flow when various volume fractions of copper-water nanofluid were present. They discovered that concentration profiles get more emaciated as Schmidt number increase. MHD convective NF flow over a rotating inclined thin layer of iron oxide based on sodium alginate was discussed by Dawar et al. [8]. The effects of magnetohydrodynamics, heat source/sink, and gravity modulation on the movement of a micropolar hybrid NF fluid across an inclined surface were studied by Ali et al. [9]. They found that fluid velocity amplifies with rising Hartmann number but emaciates with growth of micropolar material amount. The works on the related areas were seen in Refs. [10–34].

Fluid flow due to rotating conical geometries has gained importance in the field of aeronautical engineering, cooling nuclear reactor system, solar collectors, and medical field and in development of electronic chips etc. The investigation on free convection (Hering and Grosh [35]), mixed convection (Yih [36]), variable viscosity (Hossain et al. [37]), MHD (Chamkha and Al-Mudhaf [38]), and Soret/Dufour effect (Huang [39]) of nanofluids over rotating cone. The flow of MHD Ellis nanofluids around a spinning cone was perused by Ahmed et al. [40]. They stated that as bioconvection Rayleigh number ups, the temperature and velocity gradients decrease. The flow of water-copper NF and enhancement of HT between vertical concentric cones were studied by Bairi [41]. Their findings demonstrate that heat transmission improves with increasing thermal conductivity ratios, concentrations of nanoparticles, and Rayleigh numbers. The double-diffusive magnetized Jeffrey fluid flow via the disc and cone region was analysed by Khan et al. [42]. They discovered that when the thermic relaxation factor and Prandtl number increase, the thermal distribution slows down. The magnetic Casson liquid considering the impacts of Dufour and Soret because of a revolving cone was analysed by Mustafa et al. [43]. The magnetic parameter reduces the tangential and azimuthal velocities. Natural convective dissipative NF flow via a vertical cone was observed by Ragulkumar et al. [44]. They discovered that the fluid velocity and temperature of NF rise in tandem with increment in viscous dissipation parameter. The works on the related areas were carried out in Refs. [45–48].

The studied literature revealed that many works have carried out regarding the flow and HT of regular fluids or NFs over rotating cone with magnetic field. However, none of them investigated the influences of space and temperature varying heat source/sink on the stream and HT of NFs over rotating cone. The above gap in the existing literature and the relevance of NF flow over conical geometries of industrial importance were the main motivation behind the present investigation. Therefore, the objective of the existing work is to investigate the free convection flow about a vertical spinning cone under the effect of magnetic field, space and temperature varying heat source or sink using water-EG (50:50) mixture as base fluid and  $Al_2O_3$  and  $Fe_3O_4$  as nanoparticles. The novelties of the present study are:

- Introduction of space and temperature varying heat sink or source.
- Consideration of  $Al_2O_3/Fe_3O_4$  + water-EG (50:50) NF as the flow medium over a vertical spinning cone.

#### 2. Modeling

Consider an incompressible, steady, laminar and 2D magnetohydrodynamic flow of  $Al_2O_3/Fe_3O_4 + H_2O-C_2H_6O_2$  (50:50) NF along a rotating down-pointing upright hot cone subjected to magnetic field and space conjugate temperature varying heat source or sink. Cone has angular velocity shown by  $\Omega$  and the flow's local radius is taken as  $r = xsin\gamma$ . Consider u,v,w are respectively velocity elements in x,y and $\theta$  directions. Problem geometry is shown in Fig. 1. Thermophysical properties of  $H_2O-C_2H_6O_2$  (50:50) and  $Al_2O_3/Fe_3O_4$  are mentioned in Table 1.

Using the assumptions mentioned earlier, governing equations would be ([35,38,44,47]):

$$\frac{\partial (ru)}{\partial r} + \frac{\partial (rv)}{\partial v} = 0,\tag{1}$$

$$\rho_{\rm nf} \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - \frac{w^2}{x} \right) = \mu_{nf} \frac{\partial^2 u}{\partial y^2} + (\rho \beta)_{\rm nf} g \cos \gamma \left( T - T_{\infty} \right) - \sigma_{nf} B_o^2 u, \tag{2}$$

$$\rho_{\rm nf}\left(u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + \frac{uw}{x}\right) = \mu_{\rm nf}\frac{\partial^2 w}{\partial y^2} - \sigma_{\rm nf}B_{\rm o}^2w,\tag{3}$$

$$\left(\rho C_{p}\right)_{nf}\left(u\frac{\partial T}{\partial x}+v\frac{\partial T}{\partial y}\right)=k_{nf}\frac{\partial^{2}T}{\partial y^{2}}+\underbrace{\frac{k_{nf}u}{xv_{nf}}\left[A^{*}\left(T_{w}-T_{\infty}\right)exp\left(-\frac{y}{L^{2}}\right)+B^{*}\left(T-T_{\infty}\right)\right]}_{space\ and\ temperature\ dependent\ non-uniform\ heat\ source/sink}$$
(4)

Boundary conditions ([38,47]):

$$u = L_1 \frac{\partial u}{\partial y} + L_2 \frac{\partial^2 u}{\partial y^2}, v = 0, w = r\Omega + L_3 \frac{\partial w}{\partial y} + L_4 \frac{\partial^2 w}{\partial y^2}, T = T_w \text{ at } y = 0,$$
(5)

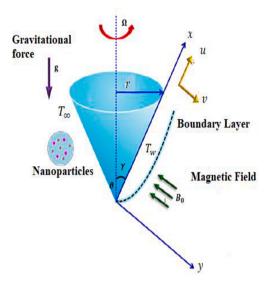


Fig. 1. Flow diagram.

Table 1 Thermal and physical properties of  $\mathbf{H}_2\mathbf{O}$ - $\mathbf{C}_2\mathbf{H}_6\mathbf{O}_2$  (50:50) base fluid and  $\mathbf{A}l_2\mathbf{O}_3/\mathbf{F}\mathbf{e}_3\mathbf{O}_4$  nanoparticles at 300K

	$C_p$	ρ	k	$\beta \times 10^{-5} \left(\frac{1}{K}\right)$	σ	Pr
EG and water (50:50) mixture $(H_20-C_2H_60_2)$ (50:50)	3288	1056	0.425	0.00341	0.00509	29.86
Aluminium oxide. $(Al_2O_3)$	765	3970	40	0.85	$35\!\times\!10^6$	-
Magnetite ( $Fe_3O_4$ )	670	5180	9.7	1.3	$2.5 \times 10^{4}$	-

$$u = 0, v = 0, T = T_{\infty} \text{ as } y \to \infty, \tag{6}$$

Thermophysical models of nanofluid are described in Table 2. Considering the transformation ([38,47]):

$$\mathbf{u} = \frac{1}{2} \Omega \mathbf{x} \sin \gamma \mathbf{F}'(\eta), \mathbf{v} = -(v\Omega \sin \gamma)^{\frac{1}{2}} \mathbf{F}(\eta), \mathbf{w} = \Omega x \sin \gamma G(\eta), = \left(\frac{\Omega \sin \gamma}{v}\right)^{1/2} \mathbf{y}, \tag{7}$$

$$\Theta(y) = \frac{T - T_{\infty}}{T_{w} - T_{\infty}} \text{ where } T_{w} - T_{\infty} = \left(T_{0} - T_{\infty}\right) \frac{x}{L},\tag{8}$$

We have

$$F''' - \left(\frac{F'^2}{2} - FF'' - 2G^2\right) (1 - \phi)^{2.5} (1 - \phi + \phi (\rho_s/\rho_f)) - \frac{\sigma_{nf}}{\sigma_f} (1 - \phi)^{2.5} MF', +2(1 - \phi)^{2.5} (1 - \phi + \phi ((\rho \beta)_s/(\rho \beta)_f)) \lambda \Theta = 0, \tag{9}$$

$$G'' - (1-\phi)^{2.5} (1-\phi + \phi(P_5/P_f)) \left(F'G - FG'\right) - \frac{\sigma_{nf}}{\sigma_f} (1-\phi)^{2.5} MG = 0, \tag{10}$$

Table 2
Thermophysical models of nanofluid.

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Properties	Nanofluid
Kinematic viscosity	$ u_{ m nf} = \frac{\mu_{ m f}}{(1-\phi)^{2.5}[(1-\phi)\rho_{ m f}+\phi\rho_{ m s}]}$
Density	$\rho_{\rm nf} = (1-\phi)\rho_{\rm f} + \phi\rho_{\rm s}$
Thermal diffusivity	$\alpha_{\rm nf} = \frac{{ m k}_{\rm nf}}{( ho{ m C}_{ m n})_{\rm of}}$
Heat capacitance	$(\rho C_p)_{nf} = (\rho C_p)_f (1 - \phi) + \phi (\rho C_p)_s$
Electrical conductivity	$\frac{\sigma_{ni}}{\sigma_i} = 1 + \frac{3(\sigma-1)\phi}{(\sigma+2)-(\sigma-1)\phi}$ , where $\sigma = \frac{\sigma_s}{\sigma_i}$
Thermal conductivity	$\frac{k_{\rm nf}}{k_{\rm f}} = \frac{(2k_{\rm f} + k_{\rm s}) - 2\phi(k_{\rm f} - k_{\rm s})}{(2k_{\rm f} + k_{\rm s}) + \phi(k_{\rm f} - k_{\rm s})}$
Thermal expansion coefficient	$(\rho\beta)_{\rm nf} = (\rho\beta)_{\rm f}(1-\phi) + \phi(\rho\beta)_{\rm s}$

$$H'' - \frac{k_f}{k_{-6}} Pr \left( 1 - \phi + \phi \left( (\rho C_p)_s J (\rho C_p)_f \right) \right) \left( \frac{1}{2} F' \Theta + \Theta' F \right) + \frac{1}{2} (1 - \phi)^{2.5} \left( 1 - \phi + \phi \left( \rho_s / \rho_f \right) \right) F' \left( A^* \exp \left( -\zeta \right) + B^* \Theta \right) = O, \tag{11}$$

With boundary conditions

$$F' = \gamma_1 F'' + \gamma_2 F''', F = 0, G = 1 + \gamma_3 G' + \gamma_4 G'', \Theta = 1 \text{ at } \eta = 0,$$
(12)

$$F' \to 0, G \to 0, \Theta \to 0 \text{ as } n \to \infty.$$
 (13)

In which  $\lambda = \frac{Gr}{Re^2}$  is mixed convection parameter,  $Re = \frac{L^2\Omega\sin\gamma}{\nu}$  denotes Reynolds number,  $Gr = \frac{g\beta_f\cos\gamma(T_0-T_\infty)L^3}{\nu^2}$  is Grashof number,  $M = \frac{\sigma B_0^2}{\rho\Omega\sin\gamma}$  is magnetic parameter,  $Pr = \frac{\nu}{\alpha}$  is Prandtl number,  $\gamma_1 = L_1\left(\frac{\Omega\sin\gamma}{\nu}\right)^{\frac{1}{2}}$  and  $\gamma_2 = L_2\left(\frac{\Omega\sin\gamma}{\nu}\right)$  are first and second order slip parameter along axial direction,  $\gamma_3 = L_3\left(\frac{\Omega\sin\gamma}{\nu}\right)^{\frac{1}{2}}$  and  $\gamma_4 = L_4\left(\frac{\Omega\sin\gamma}{\nu}\right)$  are first and order slip parameter along swirl direction.

# 3. Local nusselt number (Nu) and skin friction coefficient( $C_f$ )

Skin friction coefficients  $C_f$  and local Nusselt number Nu are defined as

$$C_{\rm f} = \frac{2\tau_{\rm w}}{\rho_{\rm f}(\Omega x \sin \gamma)^2},\tag{14}$$

and 
$$Nu = \frac{Lq_w}{k_f(T_w - T_\infty)}$$
, (15)

Where  $q_w = -k_{nf} \left(\frac{\partial T}{\partial y}\right)_{y=0}$  is surface heat flux and  $\tau_w = \mu_{nf} \left(\frac{\partial u}{\partial y}\right)_{y=0}$  indicates wall shear stress.

The dimensionless forms of  $C_{-}(f)$  & amp; Nu are

$$C_f Re^{1/2} = \left(\frac{\mu_{\rm nf}}{\mu_{\rm f}}\right) F''(0),$$
 (16)

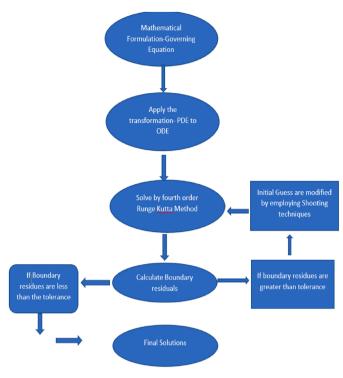
$$NuRe^{-1\backslash/2} = -\left(\frac{k_{nf}}{k_c}\right)\Theta'(0), \tag{17}$$

# 4. Numerical procedure

The BVP may well be solved via fourth order Runge Kutta scheme with shooting techniques. We set  $F = y_1$ ,  $G = y_4$ ,  $\Theta = y_6$ , and hence,

$$\begin{split} y_1' &= y_2, \\ y_2' &= y_3, \\ y_3' &= \left(\frac{F'^2}{2} - FF'' - 2G^2\right) (1 - \varphi)^{2.5} \left(1 - \varphi + \varphi \left(\rho_5/\rho_f\right)\right) + \frac{\sigma_{nf}}{\sigma_f} (1 - \varphi)^{2.5} M \ F' - 2(1 - \varphi)^{2.5} \left(1 - \varphi + \varphi \left((\rho\beta)_5/(\rho\beta)_f\right)\right) \ \lambda \ \Theta, \\ y_4' &= y_5, \\ y_5' &= (1 - \varphi)^{2.5} \left(1 - \varphi + \varphi \left(\rho_5/\rho_f\right)\right) \left(F'G - FG'\right) + \frac{\sigma_{nf}}{\sigma_f} (1 - \varphi)^{2.5} M \ G, \\ y_6' &= y_7, \\ y_7' &= \frac{k_f}{k_{nf}} \Pr \left(1 - \varphi + \varphi \left((\rho C_p)_s/(\rho C_p)_f\right)\right) \left(\frac{1}{2} F'\Theta + \Theta' F\right) - \frac{1}{2} (1 - \varphi)^{2.5} \left(1 - \varphi + \varphi \left(\rho_5/\rho_f\right)\right) F' \left(A^* \exp(-\zeta) + B^*\Theta\right). \end{split}$$

The flow chart for numerical solutions of the developed problem is provided below.



Flow chart for Numerical solutions.

#### 5. Validation

The existing outcomes are compared with that of Chamkha and Al-Mudhaf [12] and Khan et al. [26] and showed excellent agreement (See Tables 3 and 4). Ascending trend of HT rate is visualized with rise in Pr and  $\lambda$ .

# 6. Results and discussion

This part conveys the efficacies of space dependent parameter A\*, spin parameter  $\epsilon$ , magnetic parameter M, nanoparticle volume fraction  $\varphi$ , and temperature dependent parameter B\* on the temperature profile H(y), swirl G(y) and tangential F'(y) velocity profiles in the case of LST is explored. Further, variations of dimensionless  $C_f$  and Nu in response to sundry physical parameters are examined. The values used for parameters in the simulation are:  $\phi = 0.01, M = 1, \lambda = 0.3, A^* = B^* = 1, \gamma_1 = \gamma_2 = \gamma_3 = \gamma_4 = 0.01$ .

Table 3 Comparison of Skin friction coefficient (Chamkha and Al-Mudhaf [12] and Khan et al. [26]).

Pr	λ	Chamkha et al. (2005)	Khan et al. (2020)	Present result
0.7	0	1.0255	1.0256	1.020466
	1	2.2012	2.2017	2.078857
	10	8.5041	8.5044	8.524585
10	0	1.0256	1.0254	1.020465
	1	1.5636	1.5636	1.596575
	10	5.0821	5.0825	5.335744

Table 4
Comparison of Nusselt number (Chamkha and Al-Mudhaf [12] and Khan et al. [26]).

Pr	λ	Chamkha et al. (2005)	Khan et al. (2022)	Present result
0.7	0	0.4299	0.4299	0.428518
	1	0.6120	0.6120	0.612017
	10	1.0097	1.3992	1.017308
10	0	1.4110	1.4113	1.408322
	1	1.5662	1.5663	1.607681
	10	2.3580	2.3582	2.35355

#### 6.1. Influence of M on F'(v), H(v) and G(v)

Fig. 2a reveals the effect of M on F'(y) and G(y) at  $\varphi=0.1$ ,  $\varepsilon=1$ ,  $\Delta=0.1$  and Pr=29.86. When M ameliorates F'(y) and G(y) drop. This is because of the retarding nature of the Lorentz force which is due to the interaction of applied magnetic field and conducting NF thereby reduces the velocity boundary layer thickness. As visualized,  $Al_2O_3$  has higher tangential velocity compared to  $Fe_3O_4$  in M = 0 case. As M increases, magnetic nanoparticle  $Fe_3O_4$  has high tangential velocity than non-magnetic nanoparticle  $Al_2O_3$ . In case of G(y),  $Al_2O_3$  has higher spin velocity than  $Fe_3O_4$ . The influence of M on H(y) for both  $Al_2O_3$ &  $Fe_3O_4$  is portrayed in Fig. 2b. Further, H(y) ameliorates due to rise in M. However, there is subtle variation of H(y) between  $Al_2O_3$ &  $Fe_3O_4$  in appearance/non-appearance of magnetic field.

#### 6.2. Influence of $\varepsilon$ on F'(y), G(y) and H(y)

Fig. 3a depicts that increase in spin parameter  $\varepsilon$  accounts for increase in F'(y) and decrease in G(y). In fact, increase in spin parameter intensifies the movement of the nanoparticles within the NF thereby enhances the fluid motion. It is seen that Fe<sub>3</sub>O<sub>4</sub> has higher tangential velocity than  $Al_2O_3$ . In Fig. 3b it is evident that as spin parameter rises H(y) peters out and  $Al_2O_3$  has high temperature than  $Fe_3O_4$ . This implicates that heat transfer rate ameliorates for enhancement of spin parameter. In fact, rise in spin parameter enables the nanoparticle move faster thereby spend less time near the solid boundary of the cone. As a result, temperature of the NF fluid within BL emaciates.

#### 6.3. Influence of $A^*$ and $B^*$ on H(y)

The effect of  $A^*$  on H(y) is revealed in Fig. 4a. Applying  $A^* > 0$  (heat source) to BL produces energy which helps to increase H(y) while considering heat sink  $(\hat{A}(*) < 0)$  in BL may absorb energy that ensures H(y) to drop. The nanomaterial  $Al_2O_3$  has high temperature compared to  $Fe_3O_4$ . The influence of temperature dependent heat generation or absorption on temperature profile is analysed in Fig. 4b. As  $B^*$  ameliorates, temperature profile uplifts. Further,  $Al_2O_3$  has high temperature compared to  $Fe_3O_4$ .

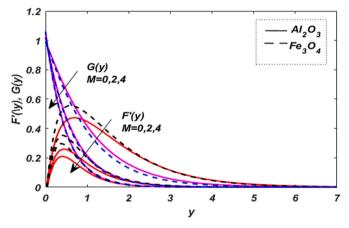


Fig. 2a. Role of M on F'(y) and G(y).

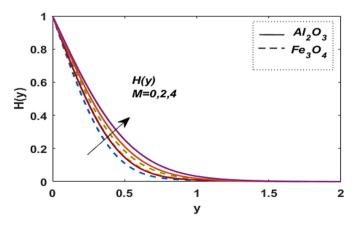


Fig. 2b. Role of M on H(y)

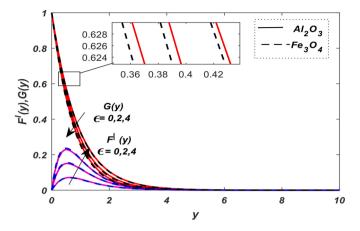
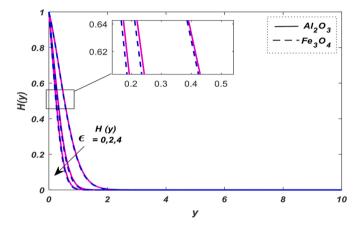


Fig. 3a. Role of  $\varepsilon$  on F'(y) and G(y).



**Fig. 3b.** Role of  $\varepsilon$  on H(y).

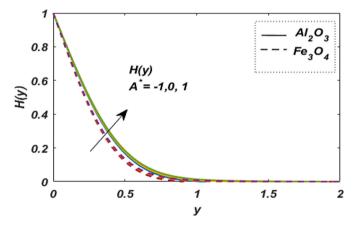


Fig. 4a. Role of  $A^*$  on H(y).

# 6.4. Influence of $\gamma_1, \gamma_2, \gamma_3, \gamma_4$ on F'(y) and H(y)

Fig. 5a, displays the impact of first order slip factor along tangential direction  $\gamma_1$  on  $F(\eta)$ . As  $\gamma_1$  increases there is enhancement in tangential velocity. The F'(y) is higher in the case of  $Fe_3O_4$  than  $Al_2O_3$ . The F'(y) profiles are characterized by overshoots near tangential surface of the cone. This is because  $F(\eta)$  is higher adjacent to solid surface than that of ambient NF. Fig. 5b indicates the effect of  $\gamma_1$  on H(y). As  $\gamma_1$  increases there is drop in temperature. Intensifying slip on the cone's surface reduces the fluid stuck to it. Consequently, fluid temperature peters out. Further,  $Al_2O_3$  has high temperature compared to  $Fe_3O_4$ . As a consequence, the NF with  $Al_2O_3$  as nanoparticle exhibits thicker thermal boundary layer. Fig. 5c displays the role of first order slip parameter along axial direction  $\gamma_2$ 

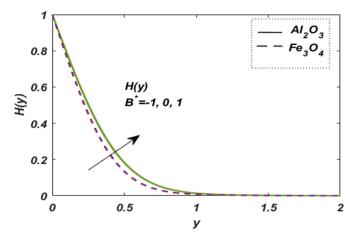
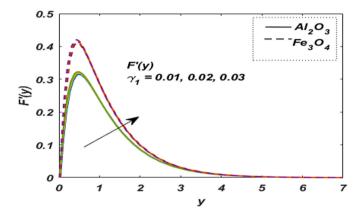
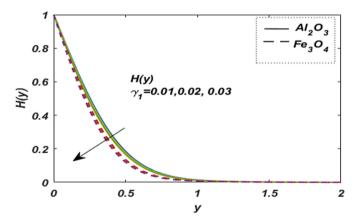


Fig. 4b. Role of  $B^*$  on H(y).

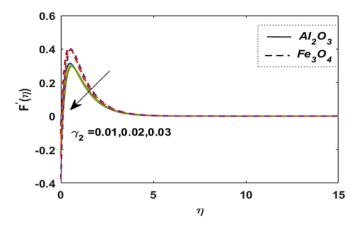


**Fig. 5a.** Role of  $\gamma_1$  on  $F'(\eta)$ .

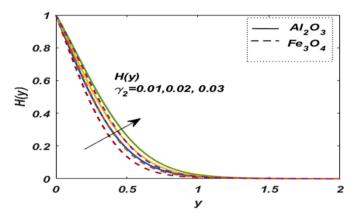


**Fig. 5b.** Role of  $\gamma_1$  on  $H(\eta)$ .

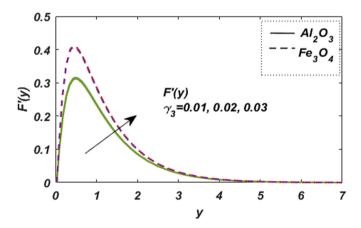
on velocity profile  $F(\eta)$ . As  $\gamma_2$  rises there is emaciation in  $F(\eta)$  and therefore, velocity boundary layer thickness decays. So rise in slip parameter  $\gamma_2$  generates decelerated flow of NFs in BL region near the surface. Amplifying slip on the cone's surface can lessen the liquid stuck to it which may lead the shear stress to diminish and consequently can impact the movement of the fluid. In addition,  $Fe_3O_4$  possesses higher velocity than that of  $Al_2O_3$ . In case of temperature profile H(y), increase in  $\gamma_2$  leads to intensification of temperature H(y). Further,  $Al_2O_3$  has high temperature than  $Fe_3O_4$  (Fig. 5d). Fig. 5e illustrates that increasing the value of  $\gamma_3$  leads to augmentation of  $F'(\eta)$ . Further,  $F'(\eta)$  of  $Al_2O_3$  nanofluid is lower than that of  $Fe_3O_4$  nanofluid. A similar trend  $\eta$  is visualized for temperature profile  $H(\eta)$  in response to rise in  $\gamma_3$  (Fig. 5f). Fig. 5g and Fig. 5h illustrate that the amplifying  $\gamma_4$  leads to enhancement of  $F'(\eta)$  and enervation in  $H(\eta)$ . In Fig. 5g,  $Al_2O_3$  is dominated by  $Fe_3O_4$ , while opposite trend is visualized in Fig. 5h.



**Fig. 5c.** Role of  $\gamma_2$  on  $F(\eta)$ .



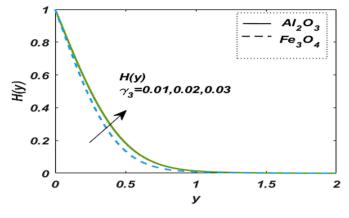
**Fig. 5d.** Role of  $\gamma_2$  on  $H(\eta)$ .



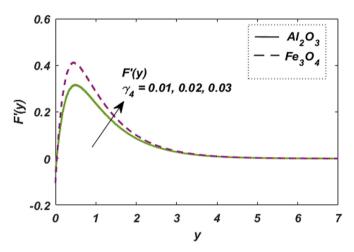
**Fig. 5e.** Role of  $\gamma_3$  on  $F(\eta)$ .

# 6.5. Influence of $\phi$ ,M.A\*,B\* on $C_fRe^{1/2}$ and NuRe<sup>-1/2</sup>

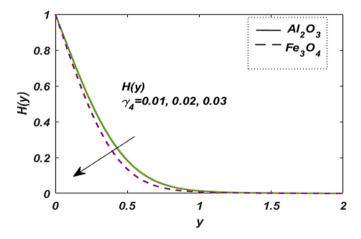
The values of  $C_fRe^{1/2}$  and NuRe<sup>-1/2</sup> for sundry  $\phi$ , M,  $A^*,B^*$  for Al<sub>2</sub>O<sub>3</sub> and Fe<sub>3</sub>O<sub>4</sub> with LST is listed in Table 5. It is noticed that NuRe<sup>-1/2</sup> decreases whereas  $C_fRe^{1/2}$  increases as  $A^*$  and  $B^*$  grow. By comparing nanoparticles  $Al_2O_3$  and  $Fe_3O_4$ ,  $Fe_3O_4$  possessessignificantly higher  $C_fRe^{1/2}$  and NuRe<sup>-1/2</sup> and NuRe<sup>-1/2</sup> and  $C_fRe^{1/2}$  peter out due to strengthening of magnetic field. However, they exhibit reverse trend with increment in volume fraction  $\phi$ .



**Fig. 5f.** Role of  $\gamma_3$  on  $H(\eta)$ .



**Fig. 5g.** Role of  $\gamma_4$  on  $F(\eta)$ .



**Fig. 5h.** Role of  $\gamma_4$  on  $H(\eta)$ .

# 7. Conclusions

The existing study deals with the free convective flow and HT of  $Al_2O_3/Fe_3O_4$  + water-EG (50:50) NF over a vertical spinning cone subject to magnetic field, space and temperature varying heat source or sink. The current problem may be utilized in the field of aeronautical engineering, cooling nuclear reactor system, solar collectors, and medical field and in development of electronic chips etc. Employing rotating conical geometries. It is worth to mention the following points from the investigation:

Table 5
The values of .C<sub>4</sub>Re<sup>1/2</sup>,and .NuRe<sup>-1/2</sup>.for sundry .φ.,M, A\*,B\*.for Al<sub>2</sub>O<sub>3</sub> and Fe<sub>3</sub>O<sub>4</sub> with linear surface temperature.

<b>φ</b> Μ		<b>A</b> *	<b>B</b> *	$C_f Re^{1/2}$		Nu <b>Re</b> - 1/2		
				$Al_2O_3$	$Fe_3O_4$	$\mathrm{Al}_2\mathrm{O}_3$	$Fe_3O_4$	
0.01	1	1	1	2.127103	3.237418	2.175797	2.398573	
0.02				3.279087	5.316236	2.445315	2.686296	
0.03				4.373248	7.292505	2.640275	2.87445	
0.01	1	1	1	2.127103	3.237418	2.175797	2.398573	
	2			1.986764	3.089628	2.03683	2.278629	
	3			1.903397	2.997533	1.942622	2.19384	
0.01	1	-2	1	2.048414	3.103143	2.235011	2.44377	
		0		2.098257	3.186411	2.196482	2.413952	
		2		2.159075	3.296437	2.153922	2.382585	
0.01	1	1	-1	2.120062	3.225694	2.195829	2.41859	
			0	2.12356	3.231517	2.185878	2.408656	
			1	2.127103	3.237418	2.175797	2.398573	
0.01	1	1	0.1	1.204577	1.641722	1.864215	2.030396	
			0.5	2.956383	4.673896	2.354564	2.584571	
			1	4.853887	7.987911	2.607726	2.808399	

- Axial velocity F(y) reduces with rise in  $M, \varepsilon, \gamma_2$  while it shows reverse trend with growth of  $\gamma_1$ .
- Swirl velocity G(y) peters out due to amplification of  $M, \varepsilon, \gamma_3$  while it reverses the trend with rise in  $\gamma_4$ .
- Temperature profiles H(y) upgrades with increase in  $M,A^*,B^*,\gamma_2,\gamma_3$  while reverse effect is attained due to rise in  $\varepsilon,\gamma_1,\gamma_4$ .
- As observed, Fe<sub>3</sub>O<sub>4</sub> exhibits higher velocity than that of Al<sub>2</sub>O<sub>3</sub> while Al<sub>2</sub>O<sub>3</sub> possesses higher temperature than that of Fe<sub>3</sub>O<sub>4</sub>.
- The surface viscous drag  $\left(C_f \operatorname{Re}^{\frac{1}{2}}\right)$  and heat transfer rate  $\left(Nu\operatorname{Re}^{-\frac{1}{2}}\right)$  enervate due to strengthening of applied magnetic field while the reverse effect is visualized with amplification of  $\phi$  for both  $Al_2O_3/Fe_3O_4$  + water nanofluids.
- 7.1. Research questions to understand the significance of the results of the current study
- What are the roles of the first and second order slip mechanisms on the behavior of velocities and temperature of magnetic nanofluid flow over rotating conical surfaces.
- Mention the influence of space and temperature varying heat generation/absorption on thermal behavior of magnetic NF flow over rotating conical surfaces.
- In what way the findings of the existing study serve scientifically to the society/industrial world?

#### CRediT authorship contribution statement

Yaqun Niu: Formal analysis, Writing. M.K. Nayak: Methodology, Writing, Simulation, Validation. S. Yashodha: Methodology, Writing, Simulation. S. Nazari: Conceptualization, Methodology, Formal analysis. A.K Abdul Hakeem: Supervision, Conceptualization. Rifaqat Ali: Writing, Formal analysis. Syed Zaheer Abbas: Formal analysis. Ali J. Chamkha: Supervision.

#### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Data availability

Data will be made available on request.

#### Nomenclature

(u, v, w) velocity elements in the x, y and  $\theta$  directions (ms<sup>-1</sup>)

x coordinate measured along the plane (m) coordinate normal to the plane (m)

MHD magnetohydrodynamic

NF nanofluid HT heat transfer BL boundary layer  $\begin{array}{ll} \rho & \quad \text{Density (Kg m}^{-3}) \\ (\rho C_p) & \quad \text{specific heat capacity} \\ \mu & \quad \text{dynamic viscosity} \\ \sigma & \quad \text{electrical conductivity} \end{array}$ 

k thermal conductivity (Wm<sup>-1</sup>K<sup>-1</sup>) β thermal expansion coefficient (K<sup>-1</sup>)

T temperature (K)

v kinematic viscosity ( $m^2s^{-1}$ )  $\gamma$  half of the vertex angle (rad)  $\Omega$  angular speed of cone ( $rads^{-1}$ ) g acceleration due to gravity ( $ms^{-2}$ )

 $B_0$  magnetic field strength

 $L_1, L_2$  first and second order slip factors along axial direction (m)  $L_3, L_4$  first and second order slip factors along swirl direction (m)

Gr Grashof number
Re Reynolds number

A\* space dependent parameterB\* temperature dependent parameter

Pr Prandtl number

λ mixed convection parameter

 $\varepsilon$  spin parameter M magnetic parameter  $C_f$  skin friction coefficient

 $\gamma_1$ ,  $\gamma_2$  first and second order slip parameters along axial direction  $\gamma_3$ ,  $\gamma_4$  first and second order slip parameters along swirl direction

Nu Nusselt number

 $\phi$  volume fraction of nanomaterial

#### Suffixes

nf nanofluidf base fluids solid nanoparticle

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